

# Reaction-diffusion equations

in biomedical applications

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  - Blood coagulation
  - Viral infection and immune response
  - Neural field models

# Reaction-diffusion equation

$u(x,t)$  – depending on considered applications describes temperature or concentration distribution, the density of some populations, etc

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

heat or mass production or reproduction rate

mass diffusion or heat conduction or random motion of individuals

One equation or systems of equations

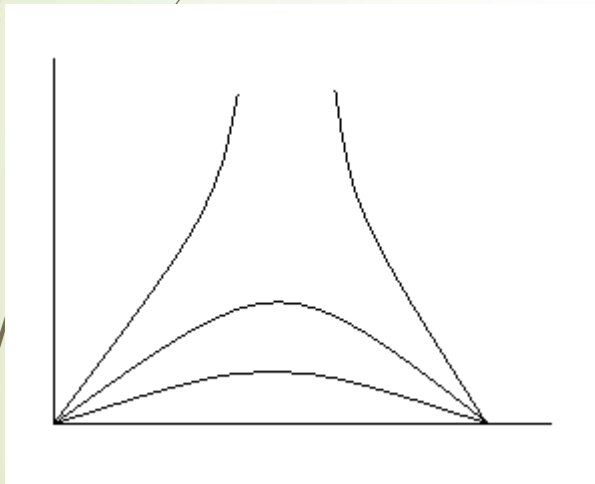
# Beginning of the theory of reaction-diffusion equations

- Heat explosion (Semenov, Frank-Kamenetskii, 1930s)
- Wave propagation (Fisher, KPP, Zeldovich-Frank-Kamenetskii, 1938)
- Pattern formation (Turing, 1952)

# Frank-Kamenetskii model of heat explosion

$$\frac{\partial u}{\partial t} = \Delta u + e^u$$

Temperature distribution



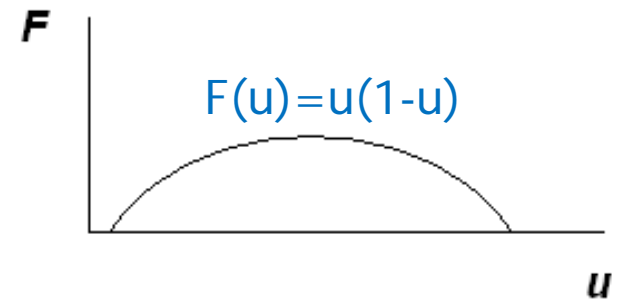
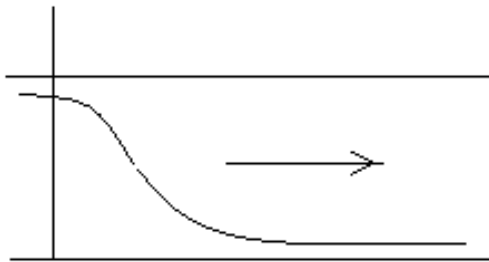
Existence of stationary solutions or blow-up (unbounded) solution?



Ammonium nitrate explosion, Beyrouth, 2020

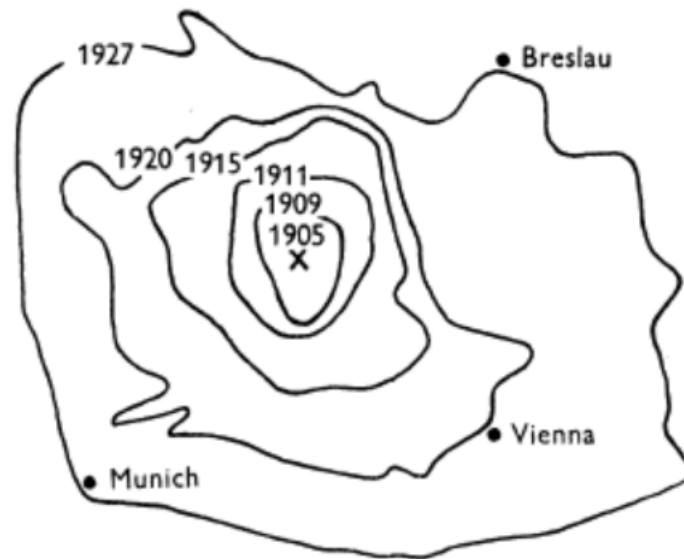
# Wave propagation (KPP-Fisher equation)

$$\frac{\partial u}{\partial t} = d \frac{\partial^2 u}{\partial x^2} + F(u)$$



$$u(x,t) = w(x-ct)$$

# Ecological invasions

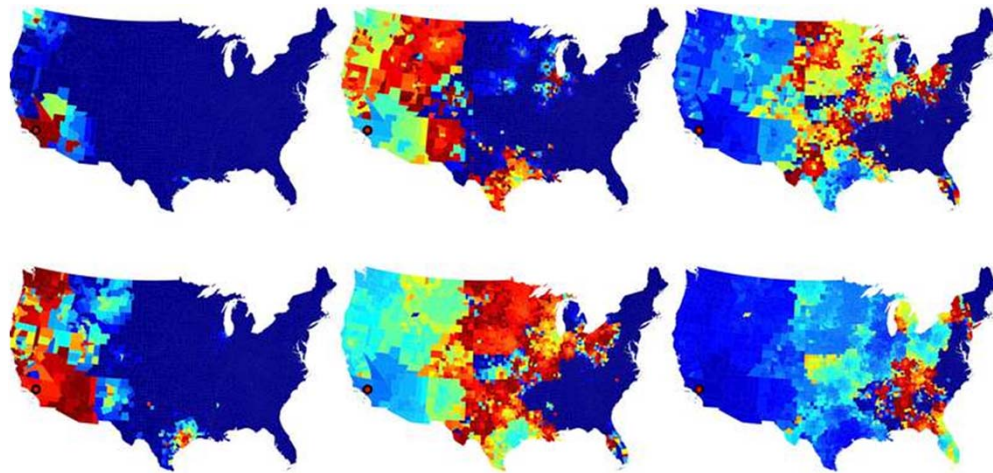


Invasion of American muskrat in Europe



Skellam, 1951

# Epidemic progression



Infection spreading in the USA

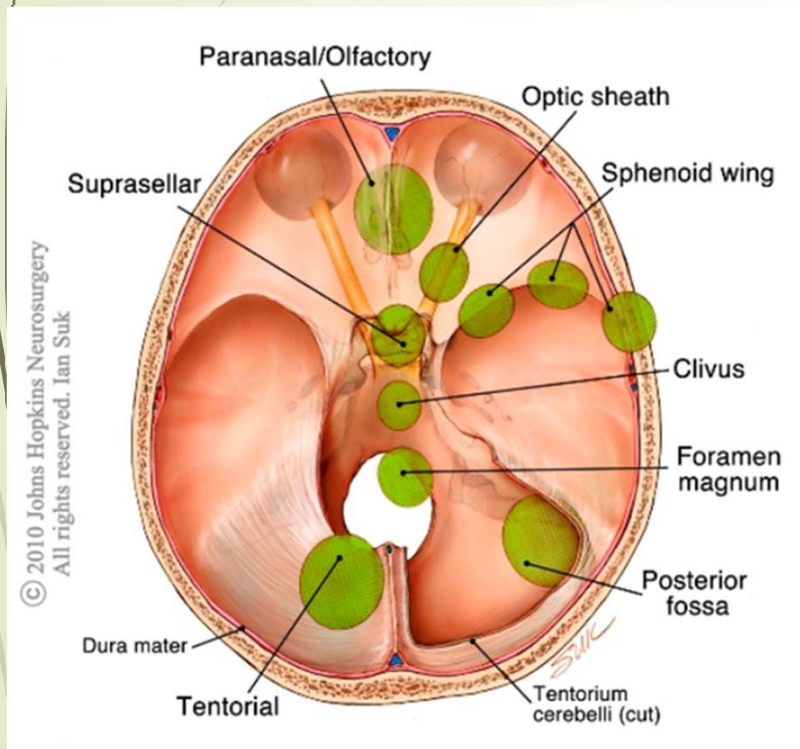


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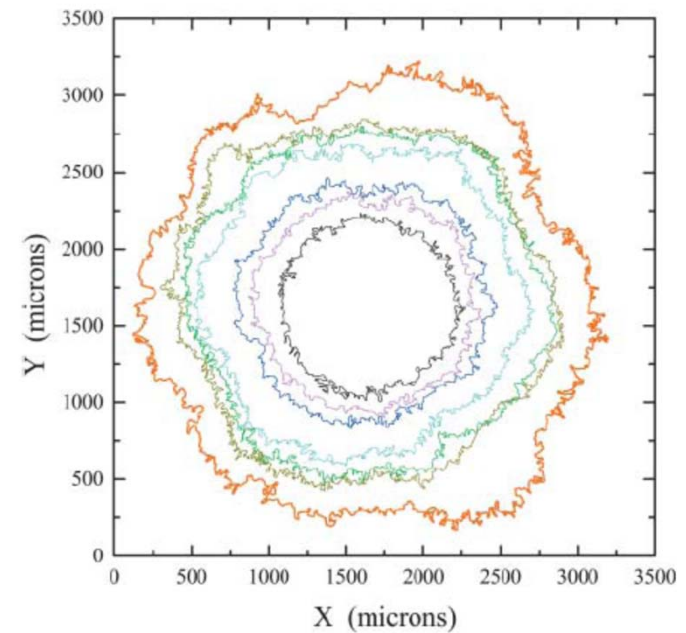
# Combustion engines, fires



# Tumor growth



$$\frac{\partial u}{\partial t} = d \frac{\partial^2 u}{\partial x^2} + au(1 - u)$$



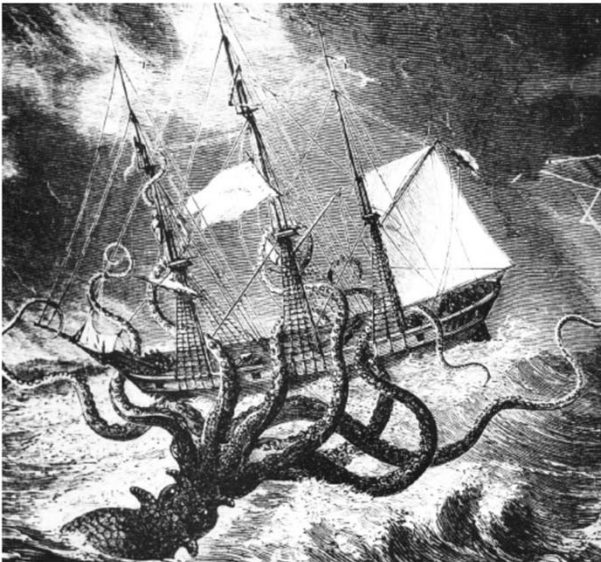
# Excitable medium: brain, heart

Nerve impulse: Hodgkin-Huxley model

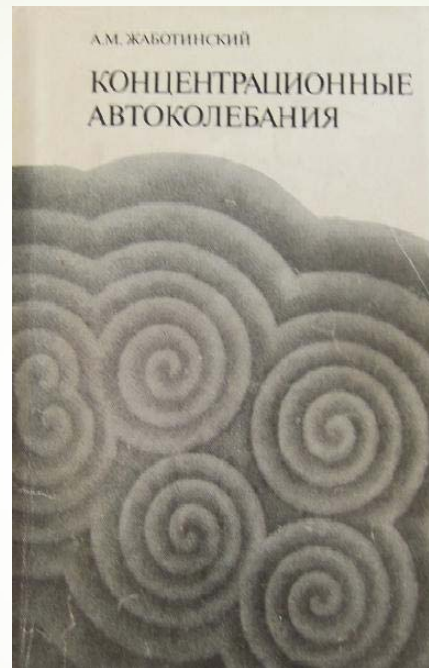
(from Keener, Sneyd)

5.1: THE HODGKIN-HUXLEY MODEL

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**Figure 5.1** The infamous giant squid (or even octopus, if you wish to be pedantic), having nothing to do with the work of Hodgkin and Huxley on squid giant axon. From *Dangerous Sea Creatures*, © 1976, 1977 Time-Life Films, Inc.

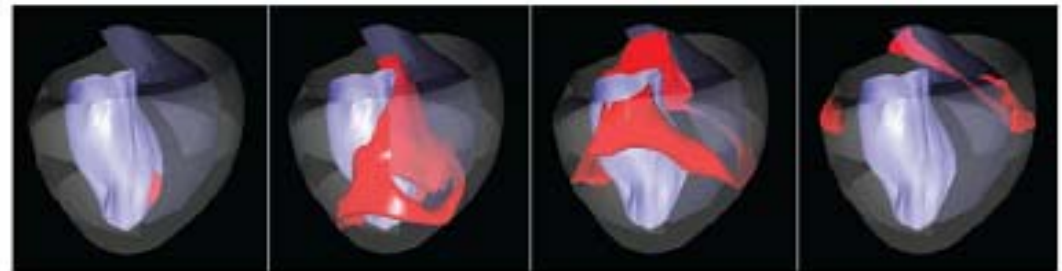


Belousov, Zhabotinskii reactions

From Zhabotinskii. Concentrational auto-oscillations. 1974.

Heart waves

(from D. Noble paper)



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## Pattern formation: Turing structures



$$\frac{\partial u}{\partial t} = d_u \frac{\partial^2 u}{\partial x^2} + F(u,v)$$

$$\frac{\partial v}{\partial t} = d_v \frac{\partial^2 v}{\partial x^2} + G(u,v)$$

# Existence and stability of waves and pulses

# Main definitions

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$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

Travelling wave solution:

$$u(x, t) = w(x - ct)$$

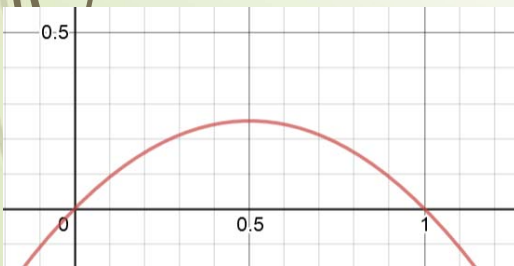
$c$  – wave speed – unknown constant

satisfies the problem:

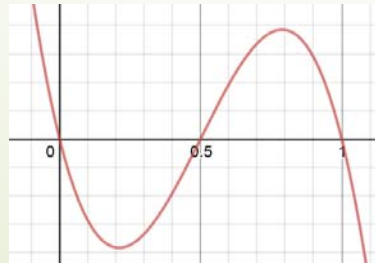
$$w'' + cw' + F(w) = 0$$

$$w(-\infty) = 1, \quad w(\infty) = 0$$

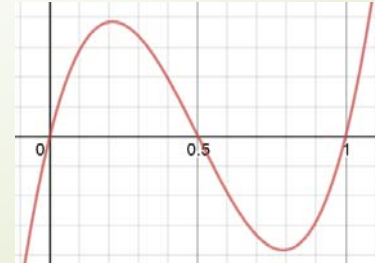
Three types of nonlinearity:



Monostable



bistable

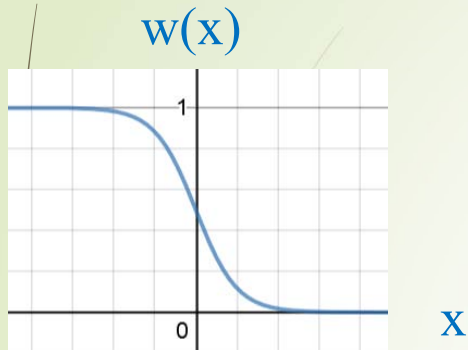


unstable

$$\frac{du}{dt} = F(u)$$

# Wave existence: monostable case

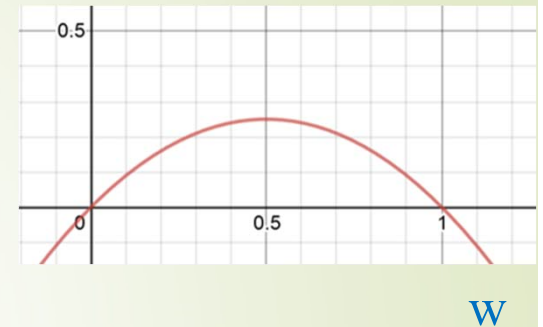
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$$w'' + cw' + F(w) = 0$$

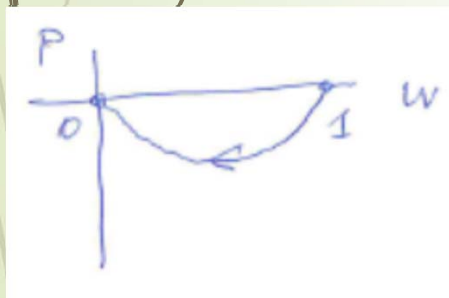
$$w(-\infty) = 1, \quad w(\infty) = 0$$

$F(w)$



First-order system:

$$w' = p, \quad p' = -cp - F(w)$$



We look for a trajectory connecting  $(1,0)$  and  $(0,0)$

## Stationary points

(0,0) Linearized system

$$\begin{cases} w' = p \\ p' = -cp - F'(0)w \end{cases} \quad \begin{pmatrix} w \\ p \end{pmatrix}' = A \begin{pmatrix} w \\ p \end{pmatrix}$$

$$A = \begin{pmatrix} 0 & 1 \\ -F'(0) & -c \end{pmatrix} \quad \boxed{F'(0) > 0}$$

Eigenvalues

$$A - \lambda E = \begin{pmatrix} -\lambda & 1 \\ -F'(0) & -c - \lambda \end{pmatrix}$$

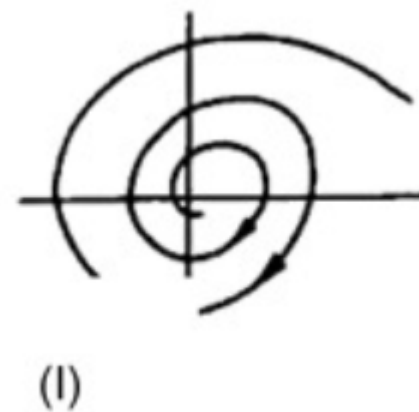
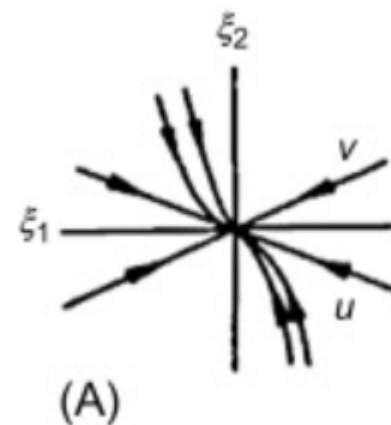
$$\det(A - \lambda E) = 0 \quad \Rightarrow$$

$$\lambda^2 + c\lambda + F'(0) = 0$$

$$\lambda = -\frac{c}{2} \pm \sqrt{\frac{c^2}{4} - F'(0)}$$

 $c > 2\sqrt{F'(0)}$  - stable node (equality?)

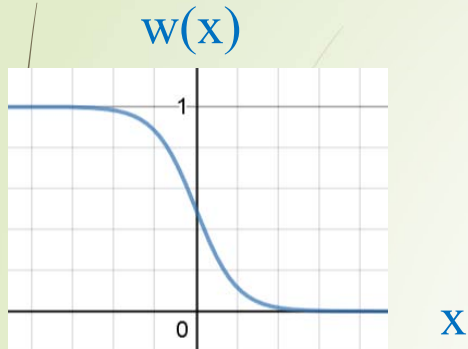
 $c < -2\sqrt{F'(0)}$  - unstable node

 $c^2 < 4F'(0)$  - focus -3-




# Wave existence: monostable case

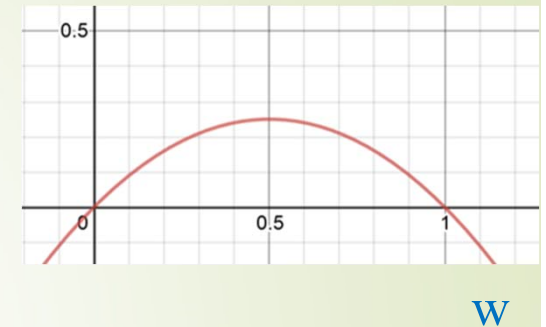
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$$w'' + cw' + F(w) = 0$$

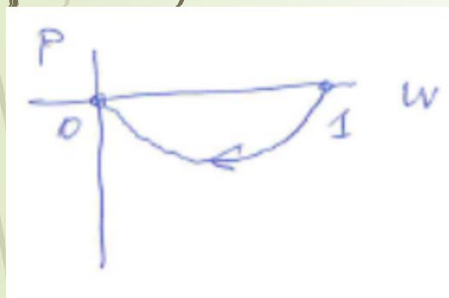
$$w(-\infty) = 1, \quad w(\infty) = 0$$

$F(w)$

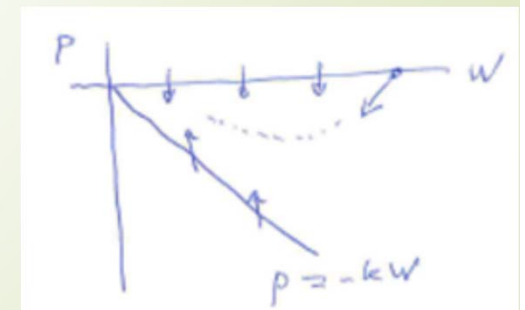


First-order system:

$$w' = p, \quad p' = -cp - F(w)$$



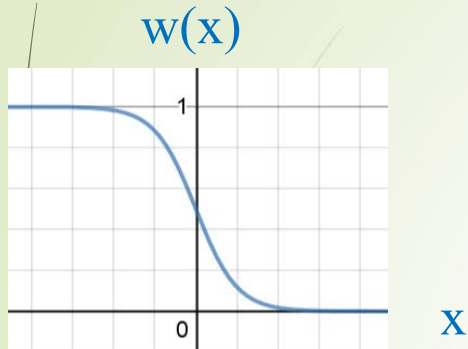
We look for a trajectory connecting  $(1,0)$  and  $(0,0)$



**Theorem 1. Monotone waves exist for all values of the speed  $c$  greater than or equal to some minimal speed  $c_0$**

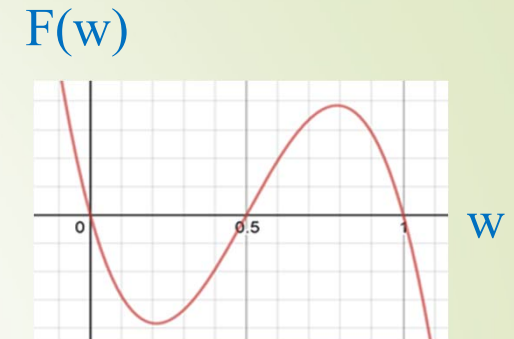
# Wave existence: bistable case

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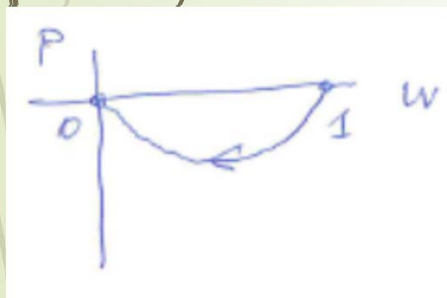
$$w'' + cw' + F(w) = 0$$

$$w(-\infty) = 1, \quad w(\infty) = 0$$

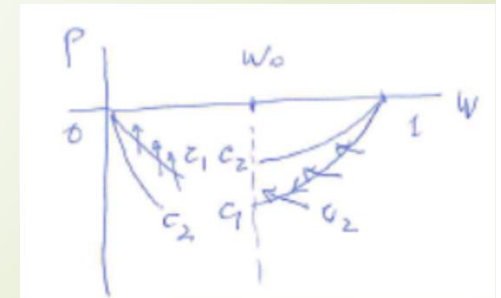


First-order system:

$$w' = p, \quad p' = -cp - F(w)$$



We look for a trajectory connecting  $(1,0)$  and  $(0,0)$



**Theorem 2.** A monotone waves exists for a single value of speed  $c$ .

# Examples

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

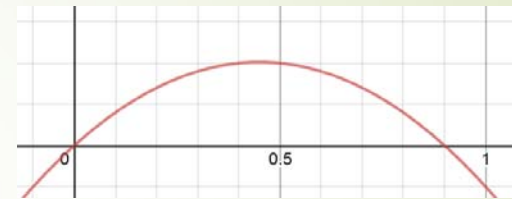
## Population dynamics

$$F(u) = au^k(1-u) - \sigma u$$

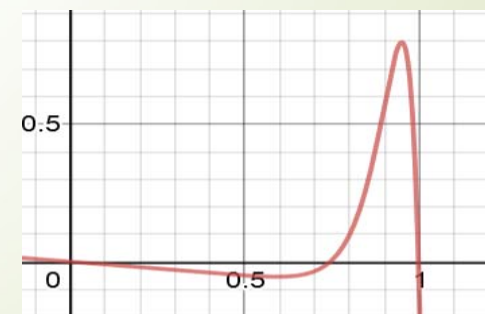
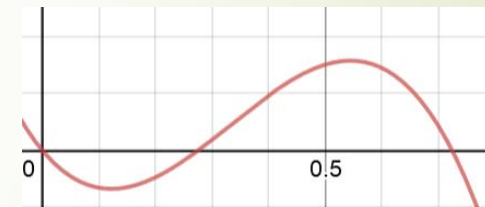
## Combustion

$$F(u) = ae^{zu}(1-u) - \sigma u$$

k=1



k=2

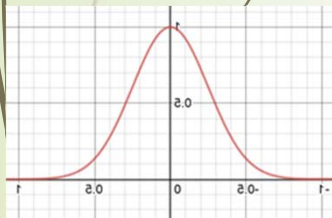


# Existence of pulses

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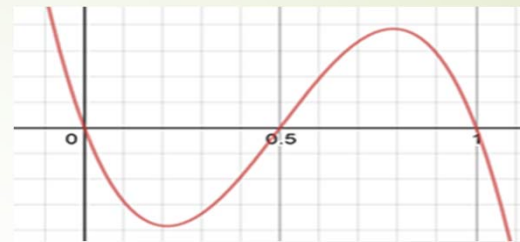
$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

Pulse – positive stationary solution with 0 limits at infinity



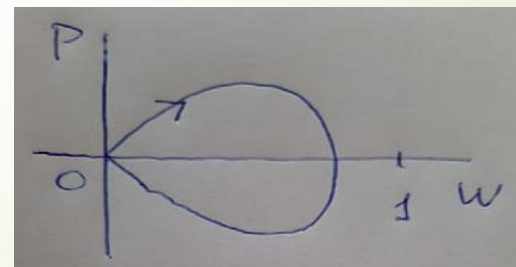
$$w'' + F(w) = 0$$

$$w(\pm\infty) = 0$$



Bistable nonlinearity

$$w' = p, \quad p' = -F(w)$$



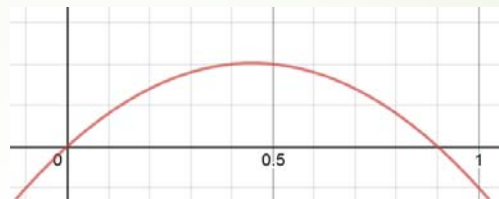
Theorem 3. Pulse exists if and only if  $\int_0^1 F(u) du > 0$

## Examples

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

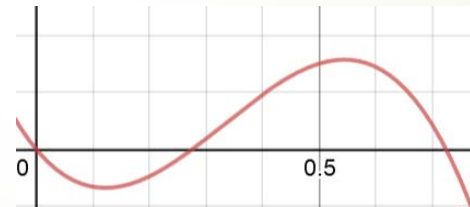
## Local model

$$F(u) = au^k(1-u) - \sigma u$$

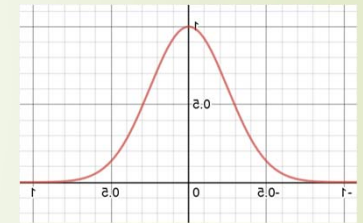


k=1

No pulse

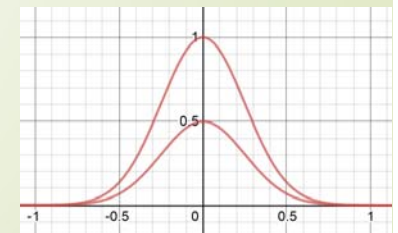
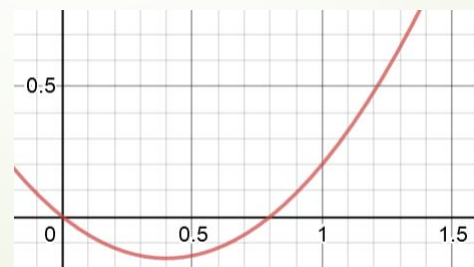


k=2



## Nonlocal model

$$F(u) = au^2 \left( 1 - \int_{-\infty}^{\infty} u(y) dy \right) - \sigma u$$



two pulses

$$r = a \left( 1 - \int_{-\infty}^{\infty} u(y) dy \right)$$

$$u'' + ru^2 - \sigma u = 0$$

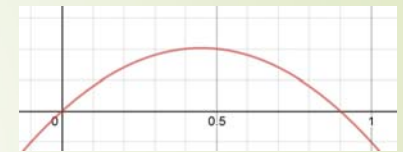
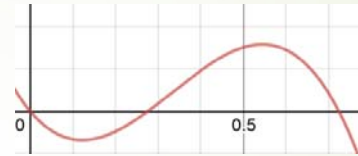
$$r = a \left( 1 - \int_{-\infty}^{\infty} u_r(y) dy \right)$$

# Generalizations

- Systems of equations
- Multi-dimensional equations and systems
- Nonlocal (integrodifferential equations)
- Delay equations

## Stability of waves: definitions

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$



$$u(x, t) = w(x - ct)$$

$$w'' + cw' + F(w) = 0$$

$$w(-\infty) = 1, \quad w(\infty) = 0$$

Invariance with respect to translation:  $w(x+h)$  is a solution for any real  $h$

## Stability of waves: definitions

Wave  $w(x)$  is a solution of the equation

$$w'' + cw' + F(w) = 0$$

Linearized operator

$$Lv = v'' + cv' + F'(w(x))v$$

Existence of zero eigenvalue:  $Lw' = 0 \rightarrow$  the previous stability result is not applicable

**Stability with shift**

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + F(u)$$

$$u(x, 0) = w(x) + \phi(x)$$

Initial condition

Asymptotic stability with shift (with respect to small perturbations): for any sufficiently small perturbation, solution  $u(x,t)$  converges to the shifted wave  $w(x-ct+h)$  for some  $h$ .



## Krein-Rutman theorem

Consider a linear elliptic operator in a bounded domains:

$$Lu = \Delta u + \sum_{i=1}^n a_i(x) \frac{\partial u}{\partial x_i} + b(x)u \quad u|_{\partial\Omega} = 0$$

The principal eigenvalue (with the maximal real part) of the problem

$$Lu = \lambda u$$

is real, simple, and the corresponding eigenfunction is positive.  
There are no other positive eigenfunctions

Remarks: valid for scalar problem and bounded domains. Can be generalized for monotone systems and unbounded domains (cf. essential spectrum)

# Stability of waves (scalar equation)

Equation for the wave  $w(x)$

(1)

$$w'' + cw' + F(w) = 0$$

Linearized  
operator

$$Lu = u'' + cu' + F'(w(x))u$$

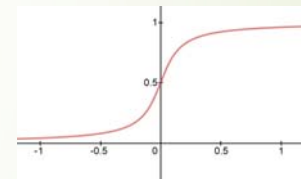
From (2): operator  $L$  has 0 eigenvalue and the corresponding eigenfunction  $w'(x)$  is positive. Hence 0 is the eigenvalue with the maximal real part, and all other eigenvalues lie in the left half-plane.

**Theorem.** Monotone waves for the scalar equation are stable.

$$u = w'(x)$$

Differentiating (1)

$$u'' + cu' + F'(w(x))u = 0 \quad (2)$$



$$u(x) = w'(x) > 0$$

# Wave speed: minimax representation

$$w'' + cw' + F(w) = 0$$



$$c = \frac{w'' + F(w)}{-w'}$$

Test function

$$\inf_x \frac{\rho'' + F(\rho)}{-\rho'} \leq c \leq \sup_x \frac{\rho'' + F(\rho)}{-\rho'}$$

Minimax representation

$$c = \inf_{\rho} \sup_x \frac{\rho'' + F(\rho)}{-\rho'} = \sup_{\rho} \inf_x \frac{\rho'' + F(\rho)}{-\rho'}$$

The proof uses global asymptotic stability of waves for monotone systems

# Instability of waves and pulses

Equation for the wave  $w(x)$

(1)

$$w'' + cw' + F(w) = 0$$

Linearized  
operator

$$Lu = u'' + cu' + F'(w(x))u$$

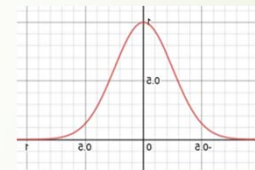
From (2): operator  $L$  has 0 eigenvalue and the corresponding eigenfunction  $w'(x)$  has variable sign. Hence 0 is not the eigenvalue with the maximal real part, and there are eigenvalues in the right half-plane.

Theorem. Non-monotone waves and pulses for the scalar equation are unstable

$$u = w'(x)$$

Differentiating (1)

$$u'' + cu' + F'(w(x))u = 0 \quad (2)$$

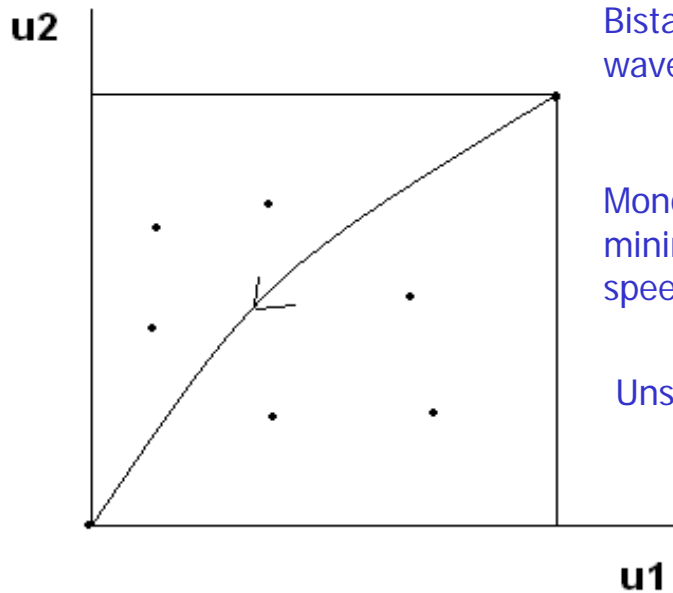


$$u(x) = w'(x)$$

# Reaction-diffusion systems: monotone systems

$$\frac{\partial u}{\partial t} = d \frac{\partial^2 u}{\partial x^2} + F(u)$$

$$\frac{\partial F_i}{\partial u_j} \geq 0, \quad i \neq j$$



Bistable: existence, uniqueness, convergence to waves, minimax representation of the speed

Monostable: existence for  $c \geq c_0$ , stability, minimax representation of the minimal wave speed

Unstable: non-existence